



#### **Phase Transformations in Materials**

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#### **Topics**

#### Single-Component Systems

- Solidification and Allotropic Transformations (Discontinuous)
  - Temperature-Dependent Thermodynamic Potentials
  - Driving Force and Latent Heat
  - Nucleation
    - Temperature Dependence
    - Homogeneous vs. Heterogeneous



#### **Temperature Dependence of** *G*





Revision of Ch. 1a



Phase Transformations in Materials

# Karlsruhe Institute of Technology





















![](_page_6_Picture_3.jpeg)

![](_page_7_Picture_0.jpeg)

![](_page_7_Figure_2.jpeg)

![](_page_7_Picture_3.jpeg)

![](_page_8_Picture_0.jpeg)

![](_page_8_Figure_2.jpeg)

![](_page_8_Picture_3.jpeg)

![](_page_9_Picture_0.jpeg)

![](_page_9_Figure_2.jpeg)

![](_page_9_Picture_3.jpeg)

![](_page_10_Picture_0.jpeg)

#### **Progress of Solidification**

![](_page_10_Figure_2.jpeg)

#### Spontaneous Solidification?

![](_page_10_Picture_4.jpeg)

#### **Progress of Solidification**

![](_page_11_Picture_1.jpeg)

![](_page_11_Figure_2.jpeg)

Metastable, super-cooled liquid!

![](_page_11_Picture_4.jpeg)

![](_page_12_Picture_1.jpeg)

![](_page_12_Figure_2.jpeg)

![](_page_12_Picture_3.jpeg)

![](_page_13_Picture_1.jpeg)

![](_page_13_Figure_2.jpeg)

![](_page_13_Picture_3.jpeg)

#### **Concept of Interface and Interphase**

![](_page_14_Picture_1.jpeg)

#### Some notes of the thermodynamic concepts on interfaces:

![](_page_14_Figure_3.jpeg)

![](_page_14_Picture_4.jpeg)

#### **Concept of Interface and Interphase**

![](_page_15_Picture_1.jpeg)

#### Some notes of the thermodynamic concepts on interfaces:

![](_page_15_Figure_3.jpeg)

![](_page_15_Picture_4.jpeg)

significantly smaller than

![](_page_15_Picture_6.jpeg)

![](_page_16_Picture_1.jpeg)

![](_page_16_Figure_2.jpeg)

![](_page_16_Picture_3.jpeg)

![](_page_17_Picture_0.jpeg)

![](_page_17_Figure_2.jpeg)

![](_page_17_Picture_3.jpeg)

![](_page_18_Picture_1.jpeg)

$$\frac{d\Delta G}{dr} = -4\pi r^2 \Delta g^{\rm D} + 8\pi r \gamma^{\rm SL}$$
$$-r^{\rm c^2} \Delta g^{\rm D} + 2r^{\rm c} \gamma^{\rm SL} = 0$$
$$r^{\rm c} = \frac{2\gamma^{\rm SL}}{\Delta g^{\rm D}}$$

$$\Delta G(r^{c}) = -\frac{4}{3}\pi \left(\frac{2\gamma^{SL}}{\Delta g^{D}}\right)^{3} \Delta g^{D} + 4\pi \left(\frac{2\gamma^{SL}}{\Delta g^{D}}\right)^{2} \gamma$$
$$= \frac{16\pi}{3}\frac{\gamma^{SL}}{\Delta g^{D}}$$

![](_page_18_Picture_4.jpeg)

![](_page_19_Picture_1.jpeg)

![](_page_19_Figure_2.jpeg)

![](_page_19_Picture_3.jpeg)

![](_page_20_Picture_1.jpeg)

![](_page_20_Figure_2.jpeg)

![](_page_20_Picture_3.jpeg)

![](_page_21_Picture_1.jpeg)

![](_page_21_Figure_2.jpeg)

![](_page_21_Picture_3.jpeg)

![](_page_22_Picture_1.jpeg)

- Apart from super cooling,  $\Delta G$  can be manipulated by modification of the interface contribution.
- For example, wetting of the nucleus at an already existing catalyst (e.g. walls of molds and dies) can be assumed:

![](_page_22_Figure_4.jpeg)

![](_page_22_Picture_5.jpeg)

![](_page_23_Picture_1.jpeg)

In order to estimate the interface energies, Young's pseudo equilibrium can be considered:

![](_page_23_Figure_3.jpeg)

Note that for mechanical equilibrium, the reaction force by the catalyst perpendicular to its surface needs to be considered by proper sectioning.

![](_page_23_Picture_5.jpeg)

## In order to estimate the interface energies, Young's pseudo equilibrium can be considered:

$$\Delta G^{\text{het}} = -V^{\text{N}} \Delta g^{\text{D}} + A^{\text{SL}} \gamma^{\text{SL}} + A^{\text{SC}} \gamma^{\text{SC}} - A^{\text{CL}} \gamma^{\text{CL}}$$
Note that the interface CL exists before nucleation and is removed by nucleation.  
Instead, the interface SC is formed.  

$$= -\frac{\pi}{3}r^{3}(2 + \cos\theta)(1 - \cos\theta)^{2} \Delta g^{\text{D}} + 2\pi r^{2} (1 - \cos\theta) \gamma^{\text{SL}} + \pi r^{2}(1 - \cos^{2}\theta)(\gamma^{\text{SC}} - \gamma^{\text{CL}})$$

$$= -\frac{\pi}{3}r^{3}(2 + \cos\theta)(1 - \cos\theta)^{2} \Delta g^{\text{D}} + 2\pi r^{2} (1 - \cos\theta) \gamma^{\text{SL}} - \pi r^{2}(1 - \cos^{2}\theta) \gamma^{\text{SL}} \cos\theta$$

$$= -\frac{\pi}{3}r^{3} \Delta g^{\text{D}} (2 - 3\cos\theta + \cos^{3}\theta) + \pi r^{2} \gamma^{\text{SL}} (2 - 3\cos\theta + \cos^{3}\theta)$$

$$= \left(-\frac{4}{3}\pi r^{3} \Delta g^{\text{D}} + 4\pi r^{2} \gamma^{\text{SL}}\right) \frac{1}{4}(2 - 3\cos\theta + \cos^{3}\theta)$$

![](_page_24_Picture_2.jpeg)

![](_page_24_Picture_6.jpeg)

![](_page_25_Picture_1.jpeg)

![](_page_25_Figure_2.jpeg)

![](_page_25_Picture_3.jpeg)

![](_page_26_Picture_1.jpeg)

![](_page_26_Figure_2.jpeg)

![](_page_26_Picture_3.jpeg)

![](_page_27_Picture_0.jpeg)

![](_page_27_Picture_2.jpeg)

critical volume V<sup>het,c</sup> decreases with better wetting

It might be counter-intuitive that the critical radius remains the same while the *G* barrier gets smaller.

The contradiction is resolved by considering the volume. Indeed, the volume of the critical nucleus gets smaller with decreasing  $\Theta$ .  $V^{\text{het,c}} = V^{\text{hom,c}} f(\Theta)$ 

The volume determines how many atoms need to join by fluctuations to form a small crystal-like object of critical size. Hence, volume determines the probability of the process in contrast to the radius of cap-shaped nucleus.

![](_page_27_Picture_7.jpeg)

#### **Heat of Fusion**

![](_page_28_Picture_1.jpeg)

![](_page_28_Picture_2.jpeg)

Note that for both melting and boiling, the entropy change  $\Delta S_{\rm m}$  or  $\Delta S_{\rm b}$  can be estimated using Trouton's rule:  $\Delta S_{\rm m} \approx \Delta S_{\rm b} \approx 10.5 k_{\rm B}$ . It is related to the change in free volume associated with the phase transition. Deviations from this rough estimate provide indication to the binding state of the respective phases.

D. R. Lide (ed.): CRC Handbook of Chemistry and Physics, Boca Raton: CRC Press LLC (2003)

![](_page_28_Picture_5.jpeg)

![](_page_29_Picture_1.jpeg)

Young's equation can be utilized by measuring (static) contact angles in images of wetted surfaces. There is a need for imaging contrast between substrate, droplet and surrounding medium:

![](_page_29_Figure_3.jpeg)

Several methods exist to recalculate the surface energy of the solid. Most commonly, the Owens, Wendt, Rabel and Kaelble (OWRK) approach is applied.

![](_page_29_Picture_5.jpeg)

![](_page_30_Picture_1.jpeg)

![](_page_30_Picture_2.jpeg)

Contact angle image glycerol on polycarbonate glass in unstrained condition. Straining the polymer causes a change in surface energy of the polymer. The images are rescaled to have equal cap radii. The scale bars are 50  $\mu$ m.

R. D. Schulman et al.: Surface energy of strained amorphous solids, Nature Communications 9 (2018) 982

![](_page_30_Picture_5.jpeg)

![](_page_31_Picture_1.jpeg)

Similar to the contact angle determination, phase/grain boundary grooves can be analyzed. In the equilibrium state of phase/grain boundary breaking through the surface, the surface tensions are in mechanical equilibrium. Diffusion from the phase/grain boundary to the free surface of the phase regions lead to the formation of grooves when being subjected to sufficiently high temperatures:

![](_page_31_Figure_3.jpeg)

![](_page_31_Picture_4.jpeg)

![](_page_32_Picture_1.jpeg)

Similar to the contact angle determination, phase/grain boundary grooves can be analyzed. The equilibrium state of phase/grain boundary penetrating the surface, the surface tensions are also in equilibrium. Diffusion from the phase/grain boundary to the free surface of phase regions lead to the formation of grooves when being subjected to sufficiently high temperatures:

![](_page_32_Figure_3.jpeg)

![](_page_32_Picture_4.jpeg)

![](_page_33_Picture_1.jpeg)

# copyright material

Atomic force microscopy (contact mode) image of grain boundary grooves at a triple junction breaking through a surface of Magnesia.

D. M. Saylor, G. S. Rohrer: Measuring the Influence of Grain-Boundary Misorientation on Thermal Groove Geometry in Ceramic Polycrystals, Journal of the American Ceramic Society 82 (1999) 1529-1536

![](_page_33_Picture_5.jpeg)

![](_page_34_Picture_1.jpeg)

Similarly, the surrounding phase can be another solid phase or the liquid in case of a three-phase equilibrium:

![](_page_34_Figure_3.jpeg)

![](_page_34_Picture_4.jpeg)

![](_page_35_Picture_1.jpeg)

Similarly, the surrounding phase can be another solid phase or the liquid in case of a three-phase equilibrium. In order to maintain the stable equ. of the liquid and the solid phases, a proper temperature gradient across the solid-liquid interface needs to be established:

![](_page_35_Figure_3.jpeg)

![](_page_35_Picture_4.jpeg)

![](_page_36_Picture_0.jpeg)

![](_page_36_Figure_2.jpeg)

Grain boundary grooves in pro-eutectic phases obtained by quenching in a Bridgman type directional solidification apparatus.

S. Engin, U. Böyük, N. Maraşlı: Determination of interfacial energies in the Al–Ag and Sn–Ag alloys by using Bridgman type solidification apparatus, Journal of Alloys and Compounds 488 (2009) 138-143

![](_page_36_Picture_5.jpeg)

![](_page_37_Picture_1.jpeg)

• For sinter contacts, reorientation of sinter particles occurs to low energy grain boundaries:

![](_page_37_Figure_3.jpeg)

![](_page_37_Picture_4.jpeg)

#### **Super Cooling**

![](_page_38_Picture_1.jpeg)

Consider Fe with  $\Delta H_{\rm m}$  = 13.81 kJ/mol,  $T_{\rm m}$  = 1811 K and a volume per atom at  $T_{\rm m}$  of  $V_{\rm A}$  = 12.7 Å<sup>3</sup> (see Ch. 2 on dilatometry). Typical values for  $\gamma^{\rm SL}$  are in the order of  $\approx 100 \text{ mJ/m}^2$ .

$$\Delta T = 20 \text{ K: } r^{\text{c}} \approx 100 \text{ Å}$$
  
 $\Delta T = 30 \text{ K: } r^{\text{c}} \approx 67 \text{ Å}$ 

This corresponds to *N*<sup>c</sup> atoms in the critical nuclei:

 $\Delta T = 10 \text{ K: } N^{\text{c}} \approx 2.7 \cdot 10^{6}$  $\Delta T = 20 \text{ K: } N^{\text{c}} \approx 3.3 \cdot 10^{5}$  $\Delta T = 30 \text{ K: } N^{\text{c}} \approx 9.9 \cdot 10^{4}$ 

![](_page_38_Picture_7.jpeg)

## Super Cooling

- Considering the high temperatures, only a very limited number of atoms need to join to a critical nucleus by fluctuations.
- In engineering systems, artificial or unintentional catalysts are always present (particles in the liquid or container walls) or perturbation of the system (convection, stirring, etc.) occurs leading to even smaller numbers of atoms to form critical nuclei, for example by heterogeneous nucleation, or to increased number of fluctuations.
- Hence, super cooling under practical conditions usually remains small.

![](_page_39_Figure_6.jpeg)

![](_page_39_Picture_7.jpeg)

#### **Super Cooling**

![](_page_40_Picture_1.jpeg)

- In order to maintain systems at large super cooling, special techniques are needed.
- There are experiments performed on International Space Station (ISS) that allow for containerless melting in low gravity conditions in order to achieve large super cooling:

![](_page_40_Picture_4.jpeg)

Dectro-static levitation furnace (ISS-ELF)

DLR (2014): https://www.dlr.de/content/de/bilder/2014/4/ohne-behaelter-schmelzen-mit-dem-eml-an-bord-der-iss-ist-das-moeglich\_17106.html NASA (2019): https://www.nasa.gov/mission\_pages/station/research/news/b4h-3rd/it-levitating-without-container

![](_page_40_Picture_7.jpeg)